



## Control design for distributed parameter systems : The port Hamiltonian approach.

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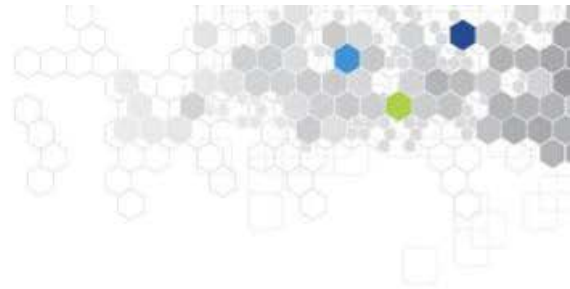
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Thank you for your attention !

## Conclusions and future works

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### Conclusion

- We provided an overview on some key results on control of distributed port Hamiltonian systems in the 1D case.
- We detailed a constructive control design technique : energy shaping for boundary/in domain controlled DPS.
- We proposed first ideas on observer design.
- We presented some possibles extensions to irreversible thermodynamic systems.

### Future works

- Study of the impact of the distribution of the patches on the achievable performances.
- Control design for a class of non linear PDE systems.
- Extension to 2D DPS.
- Control design for irreversible PHS.



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## Outline

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1. Context and motivation
2. Infinite dimensional Port Hamiltonian systems (PHS)
3. Control by interconnection and energy shaping
4. Irreversible boundary controlled port Hamiltonian Systems
5. Conclusions and future works

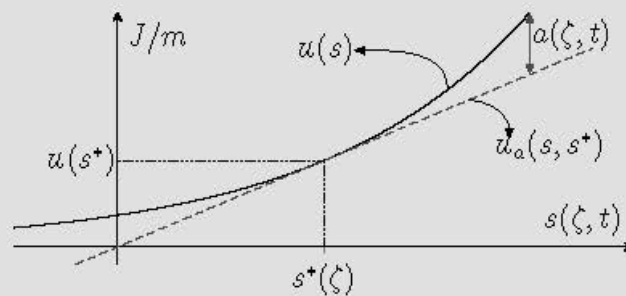


## Control design

### Idea

- Use the Thermodynamic availability function as closed loop Lyapunov function.

$$\mathcal{A} = \int_0^L (u(s) - u_a(s)) dz$$



- Use Entropy Assignment to guarantee the convergence of trajectories.

It has been successfully applied to the control of the heat equation. More complex systems (reaction-convection-diffusion systems) are under investigation.

## The non-isentropic fluid : the irreversible case

The system of balance equations may be written as the quasi-Hamiltonian system

$$\begin{bmatrix} \frac{\partial \phi}{\partial t} \\ \frac{\partial v}{\partial t} \\ \frac{\partial s}{\partial t} \end{bmatrix} = \begin{bmatrix} 0 & \frac{\partial(\cdot)}{\partial z} & 0 \\ \frac{\partial(\cdot)}{\partial z} & 0 & \frac{\partial}{\partial z} \left( \frac{\hat{\mu}}{T} \left( \frac{\partial v}{\partial z} \right) (\cdot) \right) \\ 0 & \frac{\hat{\mu}}{T} \left( \frac{\partial v}{\partial z} \right) \frac{\partial(\cdot)}{\partial z} & 0 \end{bmatrix} \begin{pmatrix} \frac{\delta H}{\delta \phi} \\ \frac{\delta H}{\delta v} \\ \frac{\delta H}{\delta s} \end{pmatrix}$$

From this new formulation (skew symmetry of the differential operator) one can define the energy/entropy boundary port variables (and input/output) such that :

$$\frac{dH}{dt} = y^T \nu$$

and

$$\frac{dS}{dt} = \underbrace{\int_a^b \sigma dz}_{\geq 0} + y_s^T \nu_s$$

## The non-isentropic fluid : the irreversible case

We can account for the **thermal domain** by considering Gibbs' equation

$$du = -pd\phi + Tds$$

where  $s$  denotes the entropy density and  $T$  the temperature. The total energy of the system is still the sum of the kinetic and the internal energy but now depends on  **$s$**

$$H(v, \phi, s) = \int_a^b \left( \frac{1}{2}v^2 + u(\phi, s) \right) dz$$

From the conservation of the total energy and Gibbs' equation  $\frac{\partial u}{\partial s} = T$  we get

$$\frac{\partial s}{\partial t}(t, z) = \frac{\hat{\mu}}{T} \left( \frac{\partial v}{\partial z} \right)^2(t, z)$$

## Irreversible systems

We consider a **1-D isentropic fluid** in Lagrangian coordinates, also known as *p-system*, with  $[a, b] \ni z$ ,  $a, b \in \mathbb{R}$ ,  $a < b$ . We choose as state variables

- the specific volume  $\phi(t, z)$ ,
- the velocity  $v(t, z)$  of the fluid.

System of two conservation laws :

$$\begin{aligned}\frac{\partial \phi}{\partial t}(t, z) &= \frac{\partial v}{\partial z}(t, z) \\ \frac{\partial v}{\partial t}(t, z) &= -\frac{\partial p}{\partial z}(t, z) - \frac{\partial \tau}{\partial z}(t, z)\end{aligned}$$

where  $p(\phi)$  is the pressure of the fluid,  $\tau = -\hat{\mu} \frac{\partial v}{\partial z}$  with  $\hat{\mu}$  the **viscous damping** coefficient. The total energy of the system is given by the sum of the kinetic energy and internal energy :

$$H(v, \phi) = \int_a^b \left( \frac{1}{2} v^2 + u(\phi) \right) dz$$

$$\begin{bmatrix} \frac{\partial \phi}{\partial t} \\ \frac{\partial v}{\partial t} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \frac{\partial}{\partial z} \left( \begin{bmatrix} \frac{\delta H}{\delta \phi} \\ \frac{\delta H}{\delta v} \end{bmatrix} \right) + \begin{bmatrix} 0 & 0 \\ 0 & \frac{\partial}{\partial z} \left( \hat{\mu} \frac{\partial}{\partial z} \right) \end{bmatrix} \left( \begin{bmatrix} \frac{\delta H}{\delta \phi} \\ \frac{\delta H}{\delta v} \end{bmatrix} \right),$$

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## Control by interconnection (Achievable performances)

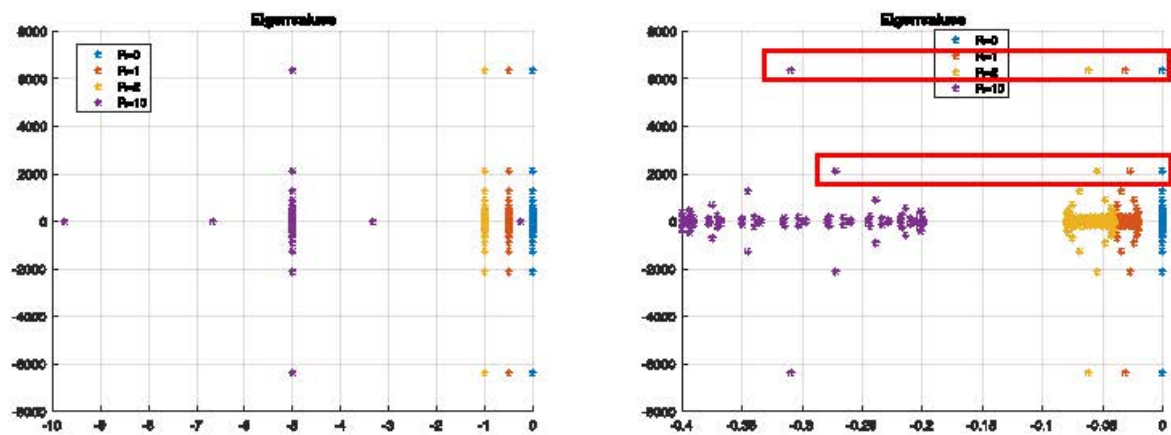
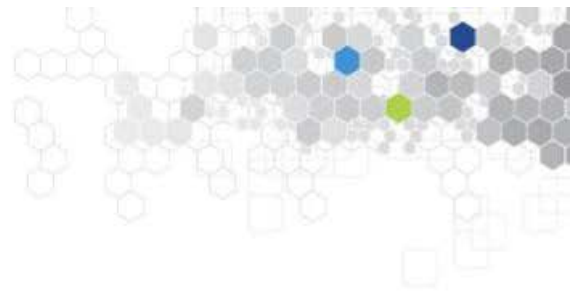


FIGURE – Control by interconnection. Full actuation (left), partial actuation (right).



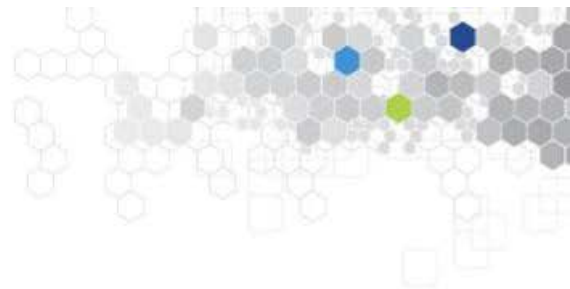
## Application case (2) (energy shaping +damping injection)

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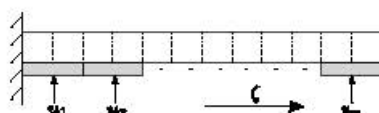
## Application case (2) (damping injection)

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## Energy shaping : application (1)

We consider the control of a weakly damped Timoshenko beam using 50 homogeneously distributed patches.



## Stability analysis

The controller is now connected to the infinite dimensional system leading to :

$$\dot{\mathcal{X}} = \underbrace{\begin{pmatrix} (\mathcal{J} - \mathcal{R} - \mathcal{B}D_c\mathcal{B}^*) & -\mathcal{B}\mathcal{B}_c^T \\ \mathcal{B}_c\mathcal{B}^* & 0 \end{pmatrix}}_{\mathcal{A}_{cl}} \begin{pmatrix} \mathcal{H} & 0 \\ 0 & Q_c \end{pmatrix} \mathcal{X}, \quad (44)$$

where  $\mathcal{X} = (x^T \quad x_c^T)^T \in X_s$  where  $X_s = L_2([0, L], \mathbb{R}^{2p}) \times \mathbb{R}^m$ .

### Existence of solution, stability analysis

- The operator  $\mathcal{A}_{cl}$  defined in (44) generates a contraction semigroup on  $X_s = L_2([0, L], \mathbb{R}^{2p}) \times \mathbb{R}^m$ .
- The operator  $\mathcal{A}_{cl}$  has a compact resolvent.
- Asymptotic stability : For any  $\mathcal{X}(0) \in L_2([0, L], \mathbb{R}^{2n}) \times \mathbb{R}^m$ , the unique solution of (44) tends to zero asymptotically, and the closed loop system (44) is globally asymptotically stable.

## Energy shaping

### Approximate energy shaping [Liu et al., 2021]

Choosing  $J_c = 0$ , and  $R_c = 0$ , the closed loop system (38) admits :

$$C(x_{1d}, x_c) = B_c M^T B_{0d}^T J_i^{-1} x_{1d} - x_c \quad (41)$$

as structural invariant along the closed loop trajectories. The control law (37) is a PI action equivalent to the state feedback :

$$u_d = -B_c^T Q_c B_c M^T B_{0d}^T J_i^{-1} x_{1d} - D_c M^T B_{0d}^T Q_2 x_{2d}. \quad (42)$$

Therefore, the closed loop system yields :

$$\begin{pmatrix} \dot{x}_{1d} \\ \dot{x}_{2d} \end{pmatrix} = \begin{pmatrix} 0 & J_i \\ -J_i^T & -(R_d + B_{0d} M D_c M^T B_{0d}^T) \end{pmatrix} \begin{pmatrix} \tilde{Q}_1 x_{1d} \\ Q_2 x_{2d} \end{pmatrix}, \quad (43)$$

where :  $\tilde{Q}_1 = Q_1 + J_i^{-T} B_{0d} M B_c^T Q_c B_c M^T B_{0d}^T J_i^{-1}$ .

$B_c^T Q_c B_c$  can be designed to minimise  $\|\tilde{Q}_1 - Q_m\|_F$  (Convex optimization problem)

## Control by interconnection

The closed loop system is given by

$$\dot{x}_{cl} = (J_{cl} - R_{cl}) Q_{cl} x_{cl}, \quad (38)$$

where  $x_{cl} = (x_{1d}^T, x_{2d}^T, x_c^T)^T$ ,  $Q_{cl} = \text{diag}(Q_1, Q_2, Q_c)$ ,

$$J_{cl} = \begin{pmatrix} 0 & J_i & 0 \\ -J_i^T & 0 & -B_{0d} M B_c^T \\ 0 & B_c M^T B_{0d}^T & J_c \end{pmatrix}, \quad R_{cl} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & R_d + B_{0d} M D_c M^T B_{0d}^T & 0 \\ 0 & 0 & R_c \end{pmatrix}.$$

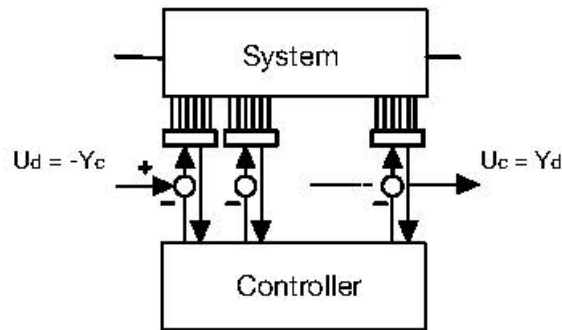
The Hamiltonian of the controller (36) is :

$$H_c(x_c) = \frac{1}{2} x_c^T Q_c x_c. \quad (39)$$

Therefore, the closed loop Hamiltonian function reads :

$$H_{cl}(x_{1d}, x_{2d}, x_c) = H_d(x_{1d}, x_{2d}) + H_c(x_c). \quad (40)$$

## Control by interconnection



The controller is designed as finite dimensional PHS of the form :

$$\begin{cases} \dot{x}_c = (J_c - R_c) Q_c x_c + B_c u_c, \\ y_c = B_c^T Q_c x_c + D_c u_c, \end{cases} \quad (36)$$

interconnected in a power preserving way through the relation

$$\begin{pmatrix} u_d \\ u_c \end{pmatrix} = \begin{pmatrix} 0 & -M \\ M^T & 0 \end{pmatrix} \begin{pmatrix} y_d \\ y_c \end{pmatrix}, \text{ where } M = \mathbb{I}_m \otimes \mathbf{1}_{k \times 1} \in \mathbb{R}^{n \times m}, \quad (37)$$



## Early lumping approach

The system is first discretized using a structure preserving method (mixed finite element method [Golo et al., 2004]) such that the approximation of (1) is again a PHS with  $n$  elements :

$$\begin{pmatrix} \dot{x}_{1d} \\ \dot{x}_{2d} \end{pmatrix} = (J_n - R_n) \begin{pmatrix} Q_1 x_{1d} \\ Q_2 x_{2d} \end{pmatrix} + B_b u_b + \begin{pmatrix} 0 \\ B_{0d} \end{pmatrix} u_d, \quad (34a)$$

$$y_b = B_b^T \begin{pmatrix} Q_1 x_{1d} \\ Q_2 x_{2d} \end{pmatrix} + D_b u_b, \quad (34b)$$

$$y_d = \begin{pmatrix} 0 & B_{0d}^T \end{pmatrix} \begin{pmatrix} Q_1 x_{1d} \\ Q_2 x_{2d} \end{pmatrix}, \quad (34c)$$

where  $x_{id} = (x_i^1 \ \dots \ x_i^n)^T \in \mathbb{R}^{np \times 1}$  for  $i \in \{1, \dots, 2p\}$ ,

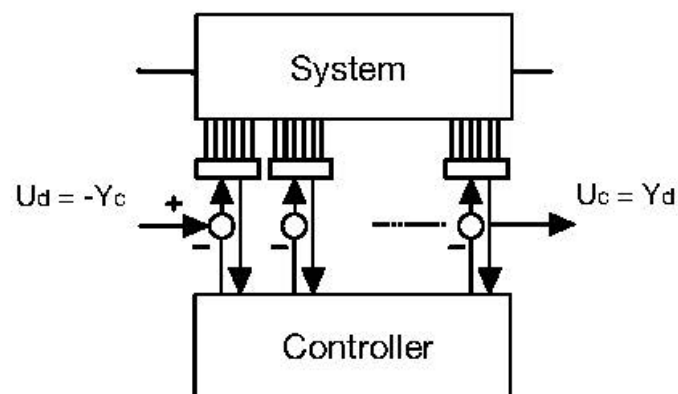
$$J_n = \begin{pmatrix} 0 & J_j \\ -J_j^T & 0 \end{pmatrix} \quad \text{and} \quad R_n = \begin{pmatrix} 0 & 0 \\ 0 & R_d \end{pmatrix},$$

The discretized energy reads :

$$H_d(x_{1d}, x_{2d}) = \frac{1}{2} \left( x_{1d}^T Q_1 x_{1d} + x_{2d}^T Q_2 x_{2d} \right). \quad (35)$$

## Control by interconnection

- **Non ideal case** : the distributed parameter system is actuated through piecewise constant elements.



## Energy shaping : ideal case

### Energy shaping [Trenchant et al., 2017]

Choosing  $\mathcal{B}_c = \mathcal{G}$  and  $\mathcal{J}_c = 0$  the closed loop system (25) admits as structural invariants the function  $C(x_\theta)$  defined by (26) and

$$\Psi = (\psi_1, 0, \psi_1)$$

In this case the hyperbolic system (1) connected to the dynamic controller (36) of the form

$$\begin{cases} \frac{\partial x_c}{\partial t}(\zeta, t) = \mathcal{G}u_c(\zeta, t) \\ y_c(\zeta, t) = \mathcal{G}^* \mathcal{Q}_c x_c(\zeta, t) + \mathcal{S}_c u_c(\zeta, t) \end{cases} \quad (31)$$

is equivalent to the system

$$\frac{\partial}{\partial t} \begin{bmatrix} x_1(\zeta, t) \\ x_2(\zeta, t) \end{bmatrix} = \begin{bmatrix} 0 & \mathcal{G} \\ -\mathcal{G}^* & -(\mathcal{R} + \mathcal{S}_c) \end{bmatrix} \begin{bmatrix} (\mathcal{H}_1(\zeta) + \mathcal{Q}_c(\zeta)) x_1(\zeta, t) \\ \mathcal{H}_2(\zeta) x_2(\zeta, t) \end{bmatrix} \quad (32)$$

$$u_\partial = \mathcal{B} \begin{bmatrix} (\mathcal{H}_1(\zeta) + \mathcal{Q}_c(\zeta)) x_1(\zeta, t) \\ \mathcal{H}_2(\zeta) x_2(\zeta, t) \end{bmatrix}, \quad y_\partial = \mathcal{C} \begin{bmatrix} (\mathcal{H}_1(\zeta) + \mathcal{Q}_c(\zeta)) x_1(\zeta, t) \\ \mathcal{H}_2(\zeta) x_2(\zeta, t) \end{bmatrix} \quad (33)$$

## Control by interconnection : ideal case

The closed loop system reads :

$$\frac{\partial x_e}{\partial t} = \begin{pmatrix} \frac{\partial x_1}{\partial t} \\ \frac{\partial x_2}{\partial t} \\ \frac{\partial x_c}{\partial t} \end{pmatrix} = \begin{pmatrix} 0 & \mathcal{G} & 0 \\ -\mathcal{G}^* & -(S_c + R) & -\mathcal{B}_c^* \\ 0 & \mathcal{B}_c & \mathcal{J}_c \end{pmatrix} \begin{pmatrix} \mathcal{H}_1 x_1 \\ \mathcal{H}_2 x_2 \\ \mathcal{Q}_c x_c \end{pmatrix} \quad (25)$$

### Structural invariants

The closed loop system (25) admits structural invariants of the form

$$\kappa_0 = C(x_e) = \int_a^b \Psi^T x_e d\zeta \quad (26)$$

with  $\Psi = (\psi_1, \psi_2, \psi_3)$  if and only if

$$-\mathcal{G}\psi_2(\zeta) = 0 = -\mathcal{B}_c\psi_2(\zeta) + \mathcal{J}_c^*\psi_3(\zeta) \quad (27)$$

$$(S_c + R)\psi_2(\zeta) = 0 \quad (28)$$

$$\mathcal{G}\psi_1(\zeta) + \mathcal{B}_c^*\psi_3(\zeta) = 0 \quad (29)$$

$$\begin{pmatrix} 0 & \mathcal{G}_1 & 0 \\ -\mathcal{G}_1^T & 0 & -\mathcal{B}_{c1} \\ 0 & \mathcal{B}_{c1}^T & \mathcal{J}_{c1} \end{pmatrix} \begin{pmatrix} \psi_1(\zeta) \\ \psi_2(\zeta) \\ \psi_3(\zeta) \end{pmatrix} \Big|_{a,b} = 0 \quad (30)$$

## Control by interconnection : ideal case

- **Ideal case** : the control acts at each point  $\zeta$  of the spatial domain.  
The controller is of the form

$$\begin{cases} \frac{\partial x_C}{\partial t}(\zeta, t) = \mathcal{J}_C \mathcal{Q}_C x_C(\zeta, t) + \mathcal{B}_C u_C(\zeta, t) \\ y_C(\zeta, t) = \mathcal{B}_C^* \mathcal{Q}_C x_C(\zeta, t) + \mathcal{S}_C u_C(\zeta, t) \end{cases} \quad (23)$$

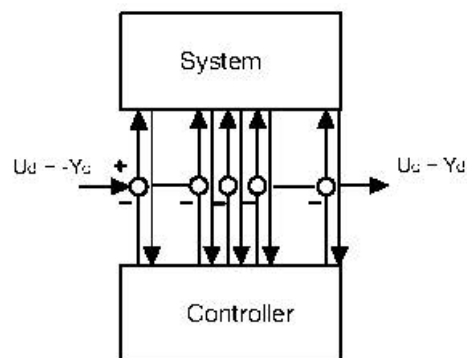
where  $\mathcal{Q}_C(\zeta) = \mathcal{Q}_C^T(\zeta)$  and  $\mathcal{Q}_C(\zeta) \geq \eta_C$  with  $\eta_C > 0$  for all  $\zeta \in [a, b]$ ,  $\mathcal{S}_C$  and  $\mathcal{S}_C(\zeta) = \mathcal{S}_C^T(\zeta)$  and  $\mathcal{S}_C(\zeta) \geq \eta_S$  with  $\eta_S > 0$  for all  $\zeta \in [a, b]$  and :

$$\mathcal{B}_C = \mathcal{B}_{C0} + \mathcal{B}_{C1} \frac{\partial}{\partial \zeta}, \text{ and } \mathcal{J}_C = \mathcal{J}_{C0} + \mathcal{J}_{C1} \frac{\partial}{\partial \zeta} \quad (24)$$

with  $\mathcal{B}_{C0}, \mathcal{B}_{C1} \in \mathbb{R}^{(n_C, 1)}$ ,  $\mathcal{J}_{C0} = -\mathcal{J}_{C0}^T$ ,  $\mathcal{J}_{C1} = \mathcal{J}_{C1}^T \in \mathbb{R}^{(n_C, n_C)}$ .

## Energy shaping

- In domain control case : we consider now in domain control



and the system is connected to the controller in a power preserving way :

$$\begin{pmatrix} u_d(\zeta, t) \\ y_d(\zeta, t) \end{pmatrix} = \begin{pmatrix} 0 & -I \\ I & 0 \end{pmatrix} \begin{pmatrix} u_c(\zeta, t) \\ y_c(\zeta, t) \end{pmatrix} + \begin{pmatrix} u'(\zeta, t) \\ 0 \end{pmatrix}, \quad (22)$$

## Implementation on the elastic string example



We consider here that

- The position of the end point *i.e.*  $\omega(b, t)$ , is measured .
- The state is reconstructed using a Luenberger PH finite dimensional observer (the controller uses the observer state) $\Rightarrow$  the closed loop stability is guaranteed [Toledo et al., 2020].



## Energy shaping

### Proposition

Under the hypothesis that the Casimir functions exist, the closed-loop dynamics (when  $u = y_c + u'$ ) is given by :

$$\begin{aligned} \frac{\partial x}{\partial t}(t, \zeta) &= P_1 \frac{\partial}{\partial \zeta} \frac{\delta H_{cl}}{\delta x}(x(t))(\zeta) + (P_0 - G_0) \frac{\delta H_{cl}}{\delta x}(x(t))(\zeta) \\ u' &= W' R \begin{pmatrix} \left( \frac{\delta H_{cl}}{\delta x}(x) \right) (b) \\ \left( \frac{\delta H_{cl}}{\delta x}(x) \right) (a) \end{pmatrix} \end{aligned} \quad (20)$$

in which  $\delta$  denotes the variational derivative, while

$$\begin{aligned} H_{cl}(x(t)) &= \frac{1}{2} \|x(t)\|_{cl}^2 + \frac{1}{2} \left( \int_a^b \hat{\Psi}^T(\zeta) x(t, \zeta) dz \right)^T \times \\ &\quad \times \hat{F}^{-1} Q_C \hat{F}^{-T} \int_a^b \hat{\Psi}(\zeta)^T x(t, \zeta) dz \end{aligned} \quad (21)$$

and  $W'$  is a  $n \times 2n$  full rank, real matrix s.t.  $W' \Sigma W'^T \geq 0$ .

## Energy shaping

- **Boundary control case** : Asymptotic stabilisation [Macchelli et al., 2017a], Exponential stabilisation [Macchelli et al., 2020]  $\Rightarrow$  Control (through  $(J_C - R_C, G_C + P_C, (G_C + P_C)^T, M_C + S_C)$ ) = integrals of the state over the spatial domain.

### Casimir functions

Consider the closed loop boundary control system with  $u' = 0$  then,

$$C(x(t), x_c(t)) = \Gamma^T x_c(t) + \int_a^b \psi^T(\zeta) x(t, \zeta) d\zeta$$

is a Casimir function for this system **if and only if**  $\psi \in H^1(a, b; \mathbb{R}^n)$ ,

$$P_1 \frac{d\psi}{d\zeta}(\zeta) + (P_0 + \mathbf{G}_0)\psi(\zeta) = 0 \quad (17)$$

$$(J_C + \mathbf{R}_C)\Gamma + (G_C + P_C)\tilde{W}R \begin{pmatrix} \psi(b) \\ \psi(a) \end{pmatrix} = 0 \quad (18)$$

$$(G_C - P_C)^T \Gamma + [W + (M_C - \mathbf{S}_C)\tilde{W}]R \begin{pmatrix} \psi(b) \\ \psi(a) \end{pmatrix} = 0 \quad (19)$$



## Energy shaping

---

- **Boundary control case** : Asymptotic stabilisation [Macchelli et al., 2017a], Exponential stabilisation [Macchelli et al., 2020]  $\Rightarrow$  Control (through  $(J_C - R_C, G_C + P_C, (G_C + P_C)^T, M_C + S_C))$  = integrals of the state over the spatial domain.

## Energy shaping

### Objectives

Modification of the closed loop system's properties (energy shaping) + stabilization (damping injection).

From the power preserving interconnection

$$H_{cl}(x, x_c) = H(x) + H_c(x_c)$$

We first look for structural invariants  $C(x, x_c)$  i.e.  $\frac{dC}{dt} = 0$

$$C(x, x_c) = x_c + F(x) = \kappa$$

where  $F$  is a smooth function. In this case the closed loop energy function reads

$$H_{cl}(x, x_c) = H_{cl}(x) = H(x) + H_c(\kappa - F(x))$$

Asymptotic stability of the closed loop system in  $x^*$  is achieved using damping injection such that

$$\frac{dH_{cl}}{dt} < 0, \forall x \neq x^*.$$

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## Energy shaping

---

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Modification of the closed loop system's properties (energy shaping) + stabilization (damping injection).



## Control by interconnection

The system is interconnected with a dynamic controller in a power preserving way.

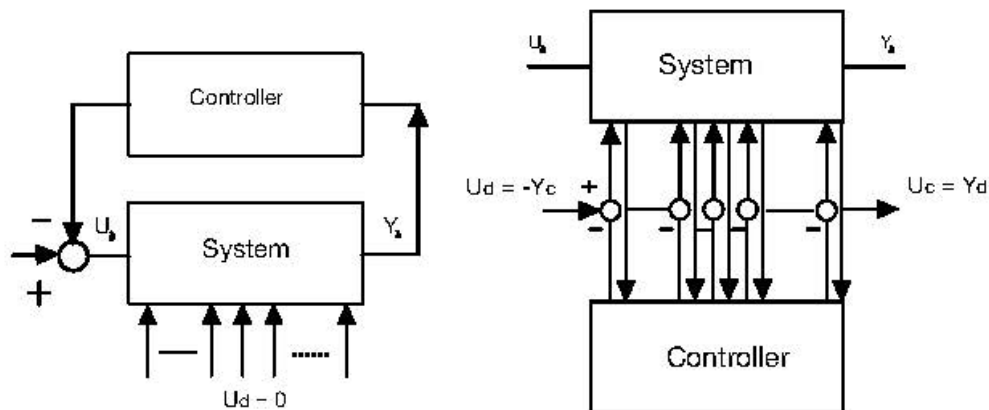
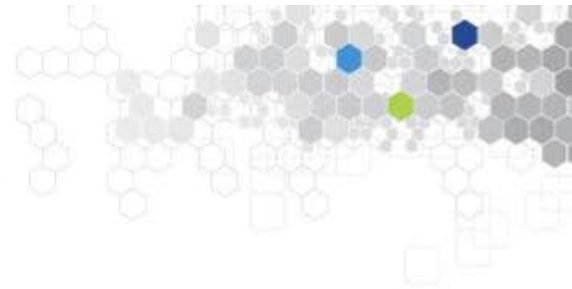


FIGURE – Control by interconnection. Boundary control (left), in domain control (right).

The closed loop energy is equal to the sum of the open loop energy and the controller energy.

## Outline

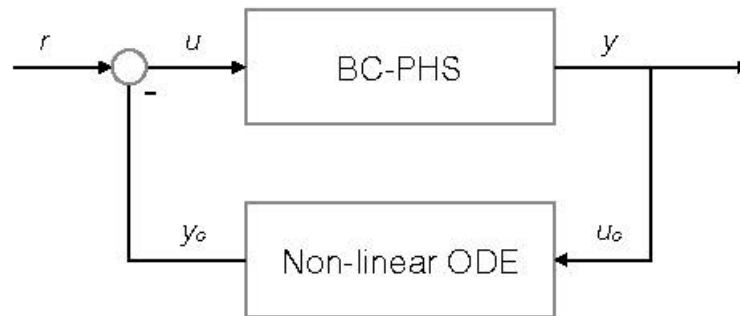
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1. Context and motivation
2. Infinite dimensional Port Hamiltonian systems (PHS)
3. Control by interconnection and energy shaping
4. Irreversible boundary controlled port Hamiltonian Systems
5. Conclusions and future works

## Non linear case

The previous results have been generalized to the non-linear case [Ramirez et al., 2017] (under some assumptions).



with

$$NL \begin{cases} \dot{v}_1 &= K_2 v_2 \\ \dot{v}_2 &= -\frac{\partial \mathcal{P}}{\partial v_1}(v_1)^T - R(K_2 v_2) + B_c u_c \\ y_c &= B_c^T K_2 v_2 + S_c u_c \end{cases} \quad (16)$$

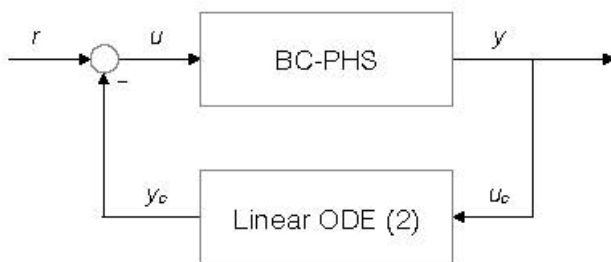
where  $v_1 \in \mathbb{R}^{n_c}$ ,  $v_2 \in \mathbb{R}^{n_c}$ , form the components of the state vector,  $B_c \in M_{k, n_c}(\mathbb{R})$ ,  $K_2 \in M_{n_c}(\mathbb{R})$ ,  $K_2 = K_2^T$ ,  $K_2 > 0$ ,  $S_c \in M_k(\mathbb{R})$  with  $S_c = S_c^T$  and  $S_c \geq 0$ .

## Dynamic boundary feedback

We consider the controller as linear finite dimensional port Hamiltonian system

$$\dot{v} = (J_c - R_c) Q_c v + B_c u_c, \quad y_c = B_c^T Q_c v + S_c u_c, \quad J_c = -J_c^T, \quad R_c = R_c^T \geq 0$$

with storage function  $E_c(t) = \frac{1}{2} \langle v(t) Q_c v(t) \rangle_{\mathbb{R}^m}$ ,  $Q_c = Q_c^T > 0 \in \mathbb{R}^m \times \mathbb{R}^m$ .



### Stability

If the following conditions are satisfied

- $\|u(t)\|^2 + \|y(t)\|^2 \geq \epsilon \|\mathcal{H}x(t, b)\|^2$ ,  $\epsilon > 0$
- power preserving interconnection  
 $u = -y_c + r$ , and  $u_c = y$
- **the controller is assumed to be exponentially stable**, i.e.,  $A_c := (J_c - R_c) Q_c$  is Hurwitz and **strictly input passive** i.e.,  $S_c > 0$ .

The closed loop system is exponentially stable.

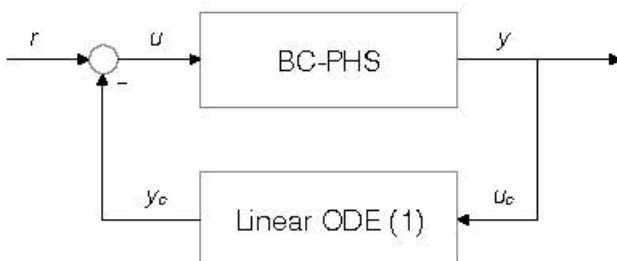
This result has been used for robust tracking control design using internal model principle [Paunonen et al., 2021].

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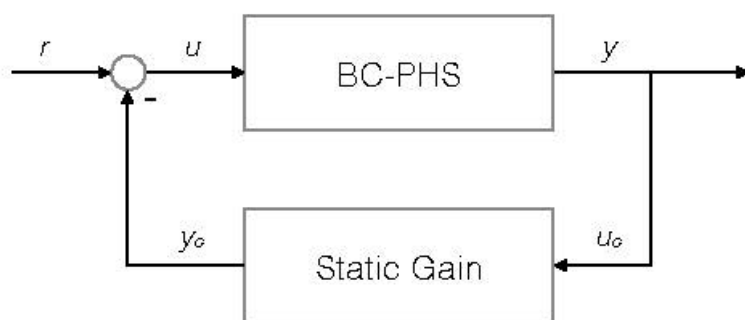
The closed loop system is asymptotically stable.

## Static feedback control

### Impedance passive case

In the impedance passive case the BCS fulfills

$$\frac{1}{2} \frac{d}{dt} \|x(t)\|_{\mathcal{H}}^2 \leq u^T(t) y(t).$$



### Static controller : $\alpha$

- Asymptotic stability :  
 $\alpha > 0$ +(compactness condition)
- Exponential stability  
[Villegas et al., 2009] :  $\alpha$  st

$$(dE/dt) \leq -k \|(\mathcal{H}x)(t, b)\|_{\mathbb{R}}^2$$

where  $k > 0$ .

This result has been used for observer design [Toledo et al., 2020].

## Static feedback control

---

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In the impedance passive case the BCS fulfills

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## The vibrating string example

The boundary port variables are

$$\begin{pmatrix} f_{\partial} \\ e_{\partial} \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} v(b) - v(a) \\ \sigma(b) - \sigma(a) \\ \sigma(b) + \sigma(a) \\ v(b) + v(a) \end{pmatrix}$$

The boundary input and output are selected as

$$u(t) = \begin{pmatrix} v(a, t) \\ \sigma(b, t) \end{pmatrix} \quad y(t) = \begin{pmatrix} -\sigma(a, t) \\ v(b, t) \end{pmatrix} \quad (15)$$

which can be derived choosing  $W$  and  $\tilde{W}$  such that :

$$W = \frac{1}{\sqrt{2}} \begin{pmatrix} -1 & 0 & 0 & 1 \\ 0 & 1 & 1 & 0 \end{pmatrix} \quad \tilde{W} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 1 & -1 & 0 \\ 1 & 0 & 0 & 1 \end{pmatrix}$$

The energy balance is then :

$$\frac{dH}{dt}(t) \leq y^T(t)u(t).$$

## The vibrating string example

The vibrating string equation is given by

$$\frac{\partial^2 \omega(\zeta, t)}{\partial t^2} = \frac{1}{\mu(\zeta)} \frac{\partial}{\partial \zeta} \left( T(\zeta) \frac{\partial \omega(\zeta, t)}{\partial \zeta} \right) - D \frac{\partial \omega(\zeta, t)}{\partial t}$$

and can be recasted in a PHS form choosing  $\varepsilon = \frac{\partial \omega(\zeta, t)}{\partial \zeta}$  and  $p = \mu \frac{\partial \omega(\zeta, t)}{\partial t}$  as state variables.

$$\frac{\partial}{\partial t} \begin{pmatrix} \varepsilon \\ p \end{pmatrix} = \begin{pmatrix} 0 & \frac{\partial}{\partial \zeta} \\ \frac{\partial}{\partial \zeta} & -D \end{pmatrix} \begin{pmatrix} T & 0 \\ 0 & \frac{1}{\mu} \end{pmatrix} \begin{pmatrix} \varepsilon \\ p \end{pmatrix}$$

which is on the form

$$\frac{\partial x}{\partial t}(\zeta, t) = \left( P_1 \frac{\partial}{\partial \zeta} + P_0 - R_0 \right) [\mathcal{H}x(\zeta, t)]$$

with

$$P_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, R_0 = \begin{pmatrix} 0 & 0 \\ 0 & D \end{pmatrix}, \mathcal{H} = \begin{pmatrix} T & 0 \\ 0 & \frac{1}{\mu} \end{pmatrix}, \mathcal{H}x(\zeta, t) = \begin{pmatrix} \sigma(\zeta, t) \\ v(\zeta, t) \end{pmatrix}$$

## Boundary controlled port Hamiltonian systems

The general formulation (1) allows to model a large class of systems.

For example :

- The 1D wave equation where  $n = 1$ ,  $N = 1$ ,  $G_0 = 0$ ,  $G_1 = 1$ .
- The Euler Bernouilli beam equation. In this case  $n = 1$ ,  $N = 2$ ,  $G_0 = 0$ ,  $G_1 = 0$ ,  $G_2 = 1$ .
- The Timoshenko beam equation. In this case  $n = 2$ ,  $N = 1$ , and

$$G_0 = \begin{bmatrix} 0 & -1 \\ 0 & 0 \end{bmatrix}, G_1 = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

In what follows we focus on first order differential operators

## Boundary controlled port Hamiltonian systems

### Existence of solution [Le Gorrec et al., 2005]

The operator

$$\mathcal{J} = \sum_{i=0}^N P_i \frac{\partial^i}{\partial \zeta^i} (\mathcal{H}(\zeta)x(\zeta, t)) - R_0 \mathcal{H}(\zeta)x(\zeta, t)$$

with domain

$$D(\mathcal{J}) = \left\{ \mathcal{H} \in H^N(a, b; \mathbb{R}^n) \mid \begin{bmatrix} \mathcal{H}(b)x(b, t) \\ \vdots \\ \frac{\partial^{N-1}(\mathcal{H}x)}{\partial \zeta^{N-1}}(b, t) \\ \mathcal{H}(a)x(a, t) \\ \vdots \\ \frac{\partial^{N-1}(\mathcal{H}x)}{\partial \zeta^{N-1}}(a, t) \end{bmatrix} \in \text{Ker} W_B \right\}$$

where  $W_B$  is defined by (9) and  $\Xi_1$  and  $\Xi_2$  satisfy (12), generates a contraction semigroup on  $X$ . Furthermore the system (5-7) with (9-10) and (12) defines a boundary control system.

## Boundary controlled port Hamiltonian systems

where

$$P_e = \begin{bmatrix} P_1 & \dots & (-1)^{N-1} P_N \\ \vdots & \ddots & 0 \\ (-1)^{N-1} P_N & 0 & 0 \end{bmatrix} \quad (11)$$

and  $\Xi_1$  and  $\Xi_2$  in  $\mathbb{R}^{k \times k}$  satisfy

$$\Xi_2^T \Xi_1 + \Xi_1^T \Xi_2 = 0, \text{ and } \Xi_2^T \Xi_2 + \Xi_1^T \Xi_1 = I \quad (12)$$

The energy balance associated to the system reads

$$\frac{dH}{dt} = \int_a^b y_d^T u_d d\zeta - \int_a^b \left( x_2^T(\zeta, t) \mathcal{H}_2^T(\zeta) R \mathcal{H}_2(\zeta) x_2(\zeta, t) \right) d\zeta + y_\partial^T u_\partial \quad (13)$$

$$\leq \int_a^b y_d^T u_d d\zeta + y_\partial^T u_\partial \quad (14)$$

## Boundary controlled port Hamiltonian systems

### Mixed in-domain / boundary controlled port Hamiltonian systems (IDBC-PHS)

A mixed in-domain / boundary controlled port Hamiltonian system is an infinite dimensional system of the form (5-7) where

$$u_{\partial} = W_B \begin{bmatrix} \mathcal{H}(b)x(b, t) \\ \vdots \\ \frac{\partial^{N-1}(\mathcal{H}x)}{\partial \zeta^{N-1}}(b, t) \\ \mathcal{H}(a)x(a, t) \\ \vdots \\ \frac{\partial^{N-1}(\mathcal{H}x)}{\partial \zeta^{N-1}}(a, t) \end{bmatrix}, \text{ and } y_{\partial} = W_C \begin{bmatrix} \mathcal{H}(b)x(b, t) \\ \vdots \\ \frac{\partial^{N-1}(\mathcal{H}x)}{\partial \zeta^{N-1}}(b, t) \\ \mathcal{H}(a)x(a, t) \\ \vdots \\ \frac{\partial^{N-1}(\mathcal{H}x)}{\partial \zeta^{N-1}}(a, t) \end{bmatrix} \quad (8)$$

with

$$W_B = \begin{bmatrix} \frac{1}{\sqrt{2}} (\Xi_2 + \Xi_1 P_{\theta}) & \frac{1}{\sqrt{2}} (\Xi_2 - \Xi_1 P_{\theta}) \end{bmatrix}, \quad (9)$$

$$W_C = \begin{bmatrix} \frac{1}{\sqrt{2}} (\Xi_1 + \Xi_2 P_{\theta}) & \frac{1}{\sqrt{2}} (\Xi_1 - \Xi_2 P_{\theta}) \end{bmatrix}, \quad (10)$$



## Infinite dimensional Port Hamiltonian systems (PHS)

For a sake of compactness we shall use the following notation

$$P_i = \begin{bmatrix} 0 & G_i \\ (-1)^{i+1} G_i^T & 0 \end{bmatrix}, \quad R_0 = \begin{bmatrix} 0 & 0 \\ 0 & R \end{bmatrix} \quad (4)$$

and the formulation of (1)

$$\frac{\partial x}{\partial t}(\zeta, t) = \sum_{i=0}^N P_i \frac{\partial^i}{\partial \zeta^i} (\mathcal{H}(\zeta)x(\zeta, t)) - R_0 \mathcal{H}(\zeta)x(\zeta, t) + \begin{bmatrix} 0 \\ I \end{bmatrix} u_d(\zeta, t) \quad (5)$$

$$y_d(\zeta, t) = \begin{bmatrix} 0 & I \end{bmatrix} \mathcal{H}(\zeta)x(\zeta, t) \quad (6)$$

$$u_\partial = \mathcal{B}(\mathcal{H}(\zeta)x(\zeta, t)), y_\partial = \mathcal{C}(\mathcal{H}(\zeta)x(\zeta, t)) \quad (7)$$

The total energy of the system  $H(x)$  is defined by

$$H(x) = \frac{1}{2} \int_a^b \left( x^T(\zeta, t) \mathcal{H}(\zeta) x(\zeta, t) \right) d\zeta$$



## Infinite dimensional Port Hamiltonian systems (PHS)

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$$\frac{\partial}{\partial t} \begin{bmatrix} x_1(\zeta, t) \\ x_2(\zeta, t) \end{bmatrix} = \begin{bmatrix} 0 & \mathcal{G} \\ -\mathcal{G}^* & -R \end{bmatrix} \begin{bmatrix} \mathcal{H}_1(\zeta)x_1(\zeta, t) \\ \mathcal{H}_2(\zeta)x_2(\zeta, t) \end{bmatrix} + \begin{bmatrix} 0 \\ I \end{bmatrix} u_d(\zeta, t) \quad (1)$$

$$y_d(\zeta, t) = \begin{bmatrix} 0 & I \end{bmatrix} \begin{bmatrix} \mathcal{H}_1(\zeta)x_1(\zeta, t) \\ \mathcal{H}_2(\zeta)x_2(\zeta, t) \end{bmatrix} \quad (2)$$

$$u_{\partial} = \mathcal{B} \begin{bmatrix} \mathcal{H}_1(\zeta)x_1(\zeta, t) \\ \mathcal{H}_2(\zeta)x_2(\zeta, t) \end{bmatrix}, \quad y_{\partial} = \mathcal{C} \begin{bmatrix} \mathcal{H}_1(\zeta)x_1(\zeta, t) \\ \mathcal{H}_2(\zeta)x_2(\zeta, t) \end{bmatrix} \quad (3)$$

where  $x = [x_1^T, x_2^T]^T \in X := L^2([a, b], \mathbb{R}^n) \times L^2([a, b], \mathbb{R}^n)$ ,  $\mathcal{H} = \text{diag}(\mathcal{H}_1, \mathcal{H}_2)$  and  $\mathcal{H}(\zeta) = \mathcal{H}^T(\zeta)$  and  $\mathcal{H}(\zeta) \geq \eta$  with  $\eta > 0$  for all  $\zeta \in [a, b]$ ,  $R \in \mathbb{R}^{(n,n)}$ ,  $R = R^T > 0$ ,  $\mathcal{B}(\cdot)$  and  $\mathcal{C}(\cdot)$  are some boundary input and boundary output mapping operators. Furthermore

$$\mathcal{G} = \sum_{i=0}^N G_i \frac{\partial^i}{\partial \zeta^i}, \quad \text{and} \quad \mathcal{G}^* = \sum_{i=0}^N (-1)^i G_i^T \frac{\partial^i}{\partial \zeta^i}$$

with  $G_i \in \mathbb{R}^{(n,n)}$ .

## Outline

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1. Context and motivation
2. Infinite dimensional Port Hamiltonian systems (PHS)
3. Control by interconnection and energy shaping
4. Irreversible boundary controlled port Hamiltonian Systems
5. Conclusions and future works

## Context : port Hamiltonian systems

---

In the linear 1D case this formalism has been used for

- Proving existence of solution using the semi-group theory [Le Gorrec et al., 2005].
- Stability analysis (when interconnected with linear or non linear ODEs) : asymptotic or exponential [Ramirez et al., 2017, Augner, 2016].
- Simulation through structure preserving schemes [Trenchant et al., 2018, Kotyczka et al., 2019].
- Control design : control by interconnection, energy shaping, observer design, backstepping ... [Macchelli et al., 2017b, Toledo et al., 2020, Redaud et al., 2022].

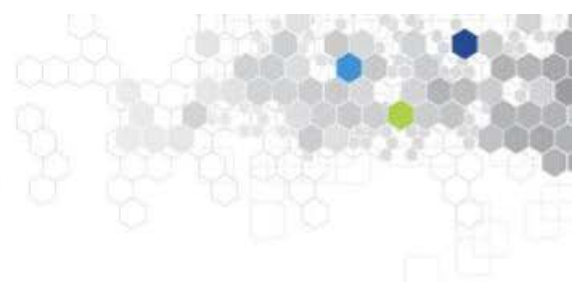
Some extensions have been proposed for

- Multidimensional systems [Skrepek, 2021].
- Implicit systems [Heidari and Zwart, 2022].
- Non linear PDE systems such as 1D or 2D-3D fluids ([Mora et al., 2021]) using Irreversible port Hamiltonian Formulations ([Ramirez et al., 2022]).

**In this talk we recall some well known results on boundary controlled port Hamiltonian systems and consider energy shaping using boundary or in domain control. Extension to IPHS.**

## Context : port Hamiltonian systems

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## Context : port Hamiltonian systems

### Port Hamiltonian systems (PHS)

Class of non linear dynamic systems derived from an **extension to open physical systems** (1992) of **Hamiltonian and Gradient systems**. This class has been generalized (2001) to distributed parameter systems.

$$x(t) : \begin{cases} \dot{x} = (J(x) - R(x)) \frac{\partial H(x)}{\partial x} + B(x)u \\ y = B(x)^T \frac{\partial H(x)}{\partial x} \\ \frac{dH}{dt} \leq y^T u \end{cases} \quad x(t, \zeta) : \begin{cases} \dot{x} = (J(x) - R(x)) \frac{\partial H(x)}{\partial x} + B_d u_d \\ y_d = B_d^T \frac{\partial H(x)}{\partial x} \\ \begin{pmatrix} f_\partial \\ e_\partial \end{pmatrix} = \frac{\partial H(x)}{\partial x} \Big|_\partial, \\ \frac{dH}{dt} \leq y_d^T u_d + f_\partial^T e_\partial \end{cases},$$

- Port Hamiltonian systems :
  - The state variables are chosen as the energy variables.
  - The links between the energy function and the system dynamics is made explicit through symmetries.
  - The boundary port variables are power conjugated.
- "Easy" to extend to non linear or systems defined on higher dimensional spaces.

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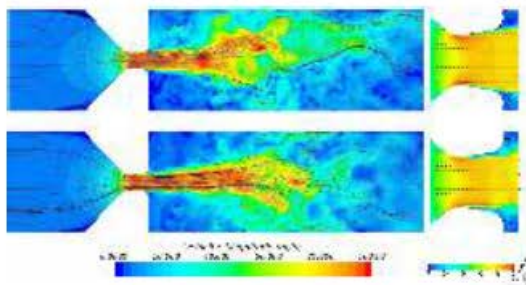
## Context / Motivation : Toward complex systems and structures

- Soft robotics (FEMTO-ST France)



- Fluid systems

- Modeling and control of interglotal air flows (coll. USM Chile)



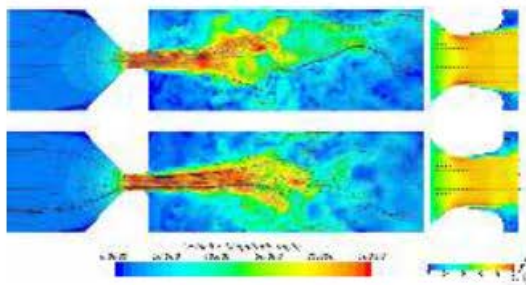
- Artificial aorta for blood pressure control (coll. EPFL Switzerland)

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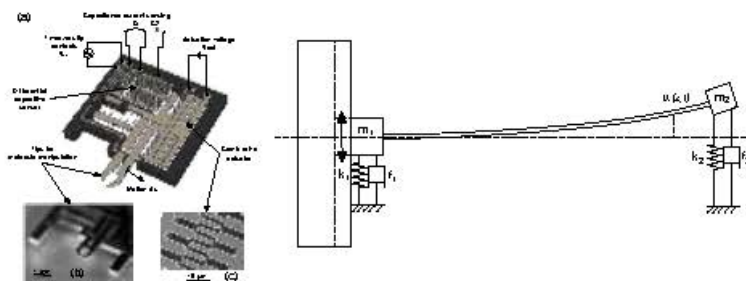
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## Context / Motivation : control of flexible structures

- Boundary controlled systems (e.g. Control of nanotweezers - Coll. LIMMS, Tokyo)



- In-domain control of distributed parameter systems (e.g. Control of smart endoscopes, FEMTO-ST)



- Exploration, imaging, diagnosis.
- Mini invasive surgery.
- Toward miniaturized and *smart* endoscopes.

## Context / Motivation : control of flexible structures

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